MoMaF : the Mock Map Facility

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ABSTRACT

We present the *Mock Map Facility*, a powerful tool to convert theoretical outputs of hierarchical galaxy formation models into catalogues of virtual observations. The general principle is straightforward : mock observing cones can be generated using semi-analytically post-processed snapshots of cosmological N-body simulations. These cones can then be projected to synthesise mock sky images. To this end, the paper describes in detail an efficient technique to create such mock cones and images from the GALICS semi-analytic model, providing the reader with an accurate quantification of the artifacts it introduces at every step. We show that replication effects introduce a negative bias on the clustering signal – typically peaking at less than 10 percent around the correlation length. We also thoroughly discuss how the clustering signal is affected by finite volume effects, and show that it vanishes at scales larger than about a tenth of the simulation box size. For the purpose of analysing our method, we show that number counts and redshift distributions obtained with GALICS/MOMAF compare well to K-band observations and to the 2dFGRS. Given finite volume effects, we also show that the model can reproduce the APM angular correlation function. The MOMAF results discussed here are made publicly available to the astronomical community through a public database. Moreover, a user-friendly Web interface (`http://galics.iap.fr`) allows any user to recover her/his own favourite galaxy samples through simple SQL queries. The flexibility of this tool should permit a variety of uses ranging from extensive comparisons between real observations and those predicted by hierarchical models of galaxy formation, to the preparation of observing strategies for deep surveys and tests of data processing pipelines.

Key words: astronomical data bases:miscellaneous - galaxies:statistics - large-scale structure of Universe - methods:numerical

1 INTRODUCTION

Large galaxy surveys in which homogeneous datasets are acquired and analysed stand to benefit from the availability of mock images of galaxy wide/deep fields. These images are very useful to design observational strategies, and to put predictions of various models into a format that can be directly compared to actual observations. Such an activity of "sky simulation" is now running extensively as part of the data processing centres for satellite missions and for the next generation of ground-based instruments, because it is now acknowledged that data processing and mock observations have to be integrated into instrument building from the beginning of the project to the interpretation of the data. However, such sky simulations are not easy, because they have to

meet a certain level of realism to be useful. The purpose of this paper is to demonstrate how *realistic* galaxy catalogues and mock images can be synthesised from the outputs of a model of hierarchical galaxy formation.

Of course, mock galaxy surveys can be generated simply by drawing galaxy types, luminosities and sizes from their distribution functions. However, such an approach does not meet the requirements, because (i) evolution cannot be included except in a very crude way, especially if we are to mimic number evolution, (ii) multi-wavelength surveys cannot be addressed easily, and (iii) spatial information cannot be addressed since galaxy positions are only known with a Poissonian distribution. Another approach is to use observed fields rescaled on-purpose (Bouwens et al. 1998). A

third approach is to start from theoretical priors and use numerical simulations to describe hierarchical clustering. In this approach, the main issue is how to transform mass to light and implement a robust method to get galaxies within dark matter structures.

There are basically two paths that this theoretical approach can take : (i) biasing schemes, and (ii) halo models. The simplest approach (i) consists of "painting" galaxies on dark matter (DM) simulations using phenomenological prescriptions to pick DM particles to be galaxies. Such methods, based on the linear bias formalism use the smoothed density field only (Cole et al. 1998) and are thus very efficient for big, low-resolution simulations (e.g. the Hubble Volume simulations described by Evrard et al. 2002). The more subtle approach (ii) uses halos identified in DM simulations. Several variations exist. In the simplest implementation, halos are populated with galaxies according to a given halo occupation distribution (Peacock & Smith 2000; Scoccimarro et al. 2001; Berlind & Weinberg 2002) which may depend on luminosity (Yang et al. 2004). In a more sophisticated approach, a Monte Carlo scheme is used to build a merger history tree for each halo identified in the simulation, and a semi-analytic model (SAM) is used to evolve galaxies in these trees (e.g. Kauffmann et al. 1997; Benson et al. 2000). In this approach, the SAM is used to generate a physically motivated halo occupation distribution. Eventually, a hybrid approach can be used, which extracts halo merging history trees from the dark matter simulations, and use a SAM to evolve galaxies in these (Kauffmann et al. 1999; Helly et al. 2003; Hatton et al. 2003, hereafter GALICS I). Mock catalogues made using this technique have been made to mimic the CfA redshift survey (Diaferio et al. 1999) or the DEEP2 survey (Coil et al. 2001). As a matter of fact, all the implementations of the two above methods are closely linked because the bias formalism often relies on analysis of the SAMs themselves (e.g. Cole et al. 1998; Somerville et al. 2001). They must therefore be considered as complementary rather than competing. The spirit of the physically motivated semi-analytic recipes can be extended to other objects such as X-ray clusters (Evrard et al. 2002). However, these approaches have generally been designed to fulfil the needs of a specific survey (for instance 2dF for Cole et al. (1998), or DEEP2 for Coil et al. (2001)). This limited scope can be extended in at least three ways.

First, it would be interesting to elaborate a generic approach to address the construction of mock observing cones from the outputs of N-body simulations at various cosmic times. The main issue is that, depending on the depth and solid angle of the observing cone, the finite size of the box may call for box replication along the line-of-sight (hereafter *radial replication*), and box replication perpendicularly to the line of sight, at the same cosmic time (hereafter *transverse replication*). Wide field, shallow surveys, with negligible evolution, can be constructed mainly from a single box (the last output of a simulation corresponding to $t = t_0$). In contrast, deep, pencil-beam surveys generally have to use numerous radial replications, whereas they may avoid transverse replication. Several issues have to be addressed here : the effects replication might have on mock catalogues, the effect using a finite volume might have on catalogues, and the sensitivity of catalogues to the number of time outputs

of the root simulation. Of course, using a larger box size would improve the situation, but given finite computer resources (CPU time and memory), using a larger box would require a trade-off in the mass resolution of the simulation, which is not acceptable if galaxies are to be modelled with a sufficient level of realism. A Hubble volume would be the ideal situation avoiding any radial or transverse replication, but so far, the largest volume simulation (ΛCDM with $3000h^{-1}$ Mpc on a side, and 10^9 particles) has a particle mass of $m_p = 2.25 \times 10^{12}h^{-1}M_\odot$ (Evrard et al. 2002), much too large to address galaxy formation with any of the "hybrid models". For instance, the ΛCDM simulation used in GALICS I has only $100h^{-1}$ Mpc on a side, and $m_p = 5.51 \times 10^9 h^{-1}M_\odot$; yet resolution effects are visible for galaxies fainter than $L^*/8$ at $z = 0$. While we await a three orders of magnitude improvement of the simulations, addressing the replication issues is unavoidable if one wants mimic large-volume observations with high resolution.

Second, the mock catalogues are useful if they gather together a large number of potentially observable properties. For instance, it is obvious that a mock catalogue designed to prepare and analyse a redshift survey of a magnitude-limited sample in a given photometric band, will incorporate at least the predicted redshifts and apparent magnitudes in that band. But the redshift survey will also be used for follow-up at other wavelengths, and other studies (for instance spectral classification once the spectra are properly calibrated). A good mock catalogue will be able to provide all these pieces of information at wavelength bands different from those of the original survey. Ultimately the mock catalogue will enable the production of field images at many wavelengths, making source extraction using the same data processing pipeline as the actual observations possible.

Third, the catalogues quickly become very large, and the question of accessibility to relevant information becomes crucial. Generally they are made available on Web pages as ASCII tables, mostly as galaxy catalogues from snapshots, more rarely as galaxy catalogues from observing cones. The more realistic these tables try to be (by including many galaxies with many properties), the more difficult to read and use they become, because of their growing size. The solution to this problem is to make the catalogues accessible through a database that can be queried to make ad hoc sub-samples fitting specific needs within a wide range of possibilities.

The purpose of this paper is to contribute along these three lines, by (i) presenting a package called MOMAF (for *Mock Map Facility*) that generates observing cones from the outputs of our GALICS model, and (ii) discussing in detail the limitations of the method. From these observing cones, synthetic catalogues are generated, that can be easily related to the catalogues of galaxies in the snapshots. The catalogues gather together a large number of properties, including magnitudes in many photometric bands of interest. These GALICS/MOMAF catalogues are made available in an on-line database that can be queried through a simple Web interface at `http://galics.iap.fr`.

To illustrate our technique, we use examples drawn from the ΛCDM simulation and the GALICS post-processing described in GALICS I. This simulation is a compromise in terms of mass resolution and volume size, and gives a satisfactory description of the luminosity functions over typically 5 mag-

nitudes. However, the techniques we describe are generic, and can be used for larger simulations. In this study, we do not address in detail the drawbacks of our model (see GALICS I) in terms of mass resolution or limited volume, nor the quality of its predictions. We are only interested in how these predictions can be converted into mock observations. First examples of using these predictions and the mock catalogues can be found in Devriendt et al., 2004 (hereafter GALICS II, in preparation), Blaizot et al. (2004, hereafter GALICS III). Two other papers will address issues which are more relevant to mock images: multi-wavelength faint galaxy counts (GALICS IV), and correlation functions (GALICS V).

This paper is organised as follows. Section 2 summarises the main features of the GALICS model that are relevant to our study. In section 3, we describe our technique of catalogue and map building from the simulation snapshots. In Section 4, we explore the different limitations of our method, most of which are actually general enough to apply to other mock catalogues in the literature based on the tiling method. We explain in section 5 how all the products of GALICS and MOMAF are stored in a relational database accessible from the web, and illustrate a few key features of this database. Section 6 contains a discussion about how these mock catalogues and images may be used, and presents perspectives for further developments.

2   THE GALICS MODEL

GALICS (for Galaxies In Cosmological Simulations) is a model of hierarchical galaxy formation which combines high resolution cosmological simulations to describe the dark matter content of the Universe with semi-analytic prescriptions to deal with the baryonic matter. This hybrid approach is fully described in GALICS I and GALICS II and we only briefly recall its relevant features here.

2.1   Dark matter simulation

The cosmological N-body simulation we refer to throughout this paper was done using the parallel tree-code developed by Ninin (1999). It is a flat cold dark matter model with a cosmological constant ($\Omega_m = 0.333$, $\Omega_\Lambda = 0.667$). The simulated volume is a cube of side $L_{box} = 100h_{100}^{-1}$Mpc, with $h_{100} = 0.667$, containing 256^3 particles of mass $8.272 \times 10^9 \mathrm{M}_\odot$, with a smoothing length of 29.29 kpc. The power spectrum was set in agreement with the present day abundance of rich clusters ($\sigma_8 = 0.88$, from Eke, Cole, & Frenk 1996), and we followed the DM density field from z=35.59 to z=0, outputting 100 snapshots spaced logarithmically in the expansion factor.

In each snapshot we use a friend-of-friend algorithm to identify virialised groups of more than 20 particles, thus setting the minimum dark matter halo mass to 1.65×10^{11} M$_\odot$. We compute a set of properties of these halos, including position and velocity of the centre of mass, kinetic and potential energies, and spin parameter. Then, assuming a density profile for the virialised dark matter, we compute the virial radius a spherical halo would have to have the same mass and potential energy, thus making the link to the idealised semi-analytic approach.

Once all the halos are identified in each snapshot, we compute their merging history trees, following the constituent particles from one output to the next one. The merging histories we obtain are by far more complex than in semi-analytic approaches as it includes evaporation of halos, fragmentation, and several artifacts due to loose friend-of-friend identifications. The way we deal with these is described in detail in GALICS I.

2.2   Baryonic Prescriptions, or how mass turns into light

When a halo is first identified, it is assigned a mass of hot gas, assuming a universal baryonic to dark matter mass ratio ($\Omega_b = 0.045$ in our fiducial model). This hot gas is assumed to be shock heated to the virial temperature of the halo, and in hydrostatic equilibrium within the dark matter potential well. The comparison of the cooling time of this gas to its free-fall time, as a function of the radius, yields the mass of gas that can cool to a central disc during a time-step. The size of this exponential disc is given by conservation of specific angular momentum during the gas in-fall and scales as the spin parameter of the halo. Then, the cooled gas is transformed into stars with a rate proportional to its mass divided by the disc dynamical time, with a given efficiency. The stars formed are distributed in mass according to an initial mass function (IMF) taken from Kennicutt (1983). The stellar population of each galaxy is then evolved between the time-steps, using a sub-stepping of at most 1 Myr. During each sub-step, stars release gas and metals in the ISM, and we follow this gas recycling in time, assuming instantaneous mixing. The massive end of the stellar population shortly explodes into supernovae which also release metals and energy in the ISM or in the IGM. We model this as a function of the instantaneous star formation rate.

When two halos merge, the galaxies they contain are gathered within the same final halo and their orbits perturbed. Subsequently, due to dynamical friction or satellite-satellite collisions, they can possibly merge. A "new" galaxy is then formed (the *descendant* of the two *progenitors*) and the stars and gas of the progenitors are distributed in three components : a disc, a bulge, and a starburst, the amount of what goes where being fixed by the ratio of masses of the two progenitors. The new galaxy can be elliptical (in shape) if the two progenitor galaxies have about the same mass, or remain a spiral if one of the merged galaxies has negligible mass.

The spectral energy distributions (SEDs) of our modelled galaxies are computed by summing the contribution of all the stars they contain, according to their age and metalicity, both of which we keep track of all along the simulation. Then, extinction is computed assuming a random inclination for disc components, and the emission of dust is added to the extinguished stellar spectra with STAR-DUST (Devriendt et al. 1999). Finally, a mean correction for absorption through the intergalactic medium (IGM) is implemented following Madau (1995), before we convolve the SEDs with the desired filters in the observer frame.

2.3   Resolution effects

The mass resolution of the DM simulations affect both the physical and statistical properties of modelled galaxies in the following ways :

(i) The particle mass of the cosmological simulation sets a minimum halo mass. Converting this halo mass into a galaxy mass, assuming that all the gas in the halo cools, one gets a threshold mass above which our sample of galaxies is complete : the *formal completeness limit*. Below this mass, although we do have galaxies, our sample is not complete since we miss galaxies in undetected halos. This direct effect of mass resolution is responsible for the lack of dwarf galaxies in the standard GALICS model. To express the completeness limit in terms of magnitudes is not straightforward because of the complex processes that convert mass into light. One can define a limiting magnitude, at a given redshift, such that, say, 95% of the galaxies brighter than that will be more massive than the formal mass resolution. Because there is no one-to-one relation between mass and luminosity, however, the luminosity selection is in practice more drastic than the selection on mass. As an example, these magnitudes are given in several wave-bands at $z = 0$ and $z = 3$, in Table 1. They can easily be derived for other wave-bands or redshifts from the GALICS database (see section 5).

(ii) In a Universe dominated by *cold* dark matter, small structures form first, and then merge and accrete material so as to evolve into larger haloes. In other words, the characteristic mass M_* of the mass distribution of haloes increases as redshift decreases. The mass resolution of our numerical simulation is fixed, however, and does not allow us to identify objects less massive than $\sim 1.6 \times 10^{11} M_\odot$, at any redshift. Hence, going back in time, more and more haloes are not resolved, and one eventually reaches a point where no halo can be detected. We call z_{lim} the limit redshift when this happens. At higher redshifts, we miss all possible galaxies. In our simulation, one find $z_{lim} \sim 7$.

(iii) A more subtle effect of resolution is that missing small structures means missing part of galaxies' histories. In practice, we showed that for our standard simulation, a galaxy needs to have evolved for about 1 Gyr before its properties have converged (see GALICS III). Although this is virtually no constraint at $z = 0$, where most galaxies are much older than 1 Gyr, the constraint becomes drastic at $z = 3$, when the age of the universe is only about 2 Gyr. To ease the selection of mature galaxies for users of the database (see section 5), we assign the morphological type 'Im' to immature galaxies.

3   MOCK OBSERVATIONS

In this section, we explain how we convert the outputs of GALICS described above into mock observations. We first describe the inputs we need from GALICS or any other model/simulation of galaxy formation. Then, we show how these inputs are turned into mock maps, and point out the main limitation of our technique : replication effects. Finally, we briefly explain how we can project catalogues onto realistic pre- or post-observing maps.

	Wave-band	95 % completeness	75 % completeness
	U	-19.6 mag	-18.2 mag
	B	-19.5 mag	-18.5 mag
	V	-20.1 mag	-19.4 mag
$z = 0$	R	-20.6 mag	-20.1 mag
	I	-21.1 mag	-20.7 mag
	J	-21.8 mag	-21.5 mag
	K	-22.8 mag	-22.5 mag
	U_n	-20.6 mag	-20.3 mag
$z = 3$	G	-21.1 mag	-20.9 mag
	R	-21.1 mag	-20.9 mag

Table 1. 95% and 75% completeness limits in terms of absolute rest-frame magnitudes at redshifts 0 and 3. At $z = 0$, the magnitudes are expressed in the Vega system, and the filters are Johnson's. At $z = 3$, the magnitudes are expressed in the AB system, and the filters are those from Steidel & Hamilton (1993).

3.1   Inputs

The method we developed to generate mock catalogues from outputs (or *snapshots*) of cosmological simulations at a finite number of redshifts is general and can be used for a variety of objects (e.g. clusters or quasars). Here, we describe the features needed in these snapshots for galaxies.

The snapshots have to be (cubic) volumes of equal comoving size (in our standard simulation, $L_{box} = 100$ h^{-1}Mpc) with periodic boundary conditions. These snapshots must each contain the following information :

• The redshift, or expansion factor of the snapshot.
• The position of each galaxy within the snapshot.
• The velocity of each galaxy within the snapshot.
• The characteristic scale-length of each component of each galaxy (disc, bulge and burst).
• The inclination of each galaxy, which was used to compute its extinction.
• The absolute AB magnitude of each galaxy in the desired filters, computed in the observer frame as

$$M_{\nu_0}(z) = -2.5 \log \left(\frac{L_{\nu_0(1+z)}}{[10 \text{ pc}]^2} \right) - 2.5 \log(1+z) + 48.6, \quad (1)$$

where z is the redshift of the snapshot, and

$$L_{\nu_0(1+z)} = \int (1+z) f_{\nu_0}(\nu) L[\nu(1+z)] \mathrm{d}\nu \qquad (2)$$

is the luminosity of a galaxy at redshift z, through a normalised filter response f_{ν_0}. Note that because the peculiar velocities or positions relative to the observer are not known at this stage, we use the redshift z of the snapshot to compute these magnitudes. This approximation will be corrected for when we compute apparent magnitudes (see Sec. 3.2.2). Also note that these magnitudes take into account extinction by the intergalactic medium, computed at the redshift of the snapshot.

• The first order derivatives of the above magnitudes with redshift, in each filter. For galaxies in snapshot i, these derivatives are estimated as

$$\frac{\mathrm{d}M}{\mathrm{d}z} = \frac{M[z(i-1)] - M[z(i)]}{z(i-1) - z(i)}, \qquad (3)$$

where $z(i)$ is the redshift of snapshot i (in our convention, $z(i) < z(i-1)$), and $M[z]$ the observer-frame absolute magnitude assuming the galaxy is at redshift z (Eq. 1). Note that this expression does not account for the evolution of galaxies, as it involves the magnitudes of the *same* galaxy put at different redshifts. Eq. 3 however captures K-correction and variations of IGM extinction with z, which are the main drivers of average variations of apparent properties with redshift in mock catalogues (see Sec. 3.2.2 and 4.3.1).

All the above quantities are direct outputs of GALICS, except for positions and velocities. These are computed as a post-treatment, using information from the DM simulations (positions and velocities of the halos) and from GALICS (orbital radii). The position of a galaxy within a snapshot is thus defined by $\vec{g} = \vec{h} + r_{orb} \times \vec{u}$, where $\vec{h}$ is the position of its host halo, r_{orb} the orbital radius of the galaxy, and $\vec{u}$ a normalised vector of random direction. The peculiar velocity of a galaxy is defined as $\vec{v}_g = \vec{v}_h + \delta\vec{v}$, where $\vec{v}_h$ is the peculiar velocity of its host halo, and $\delta\vec{v}$ is the peculiar velocity of the galaxy within this halo. The amplitude of $\delta\vec{v}$ is drawn randomly from a Gaussian distribution of width equal to the circular velocity of the halo, and its direction is random. Note that the velocities of central galaxies are taken to be that of the centre of mass of their host haloes.

3.2 Mock Catalogues

Such inputs, corresponding to the same simulated region of universe at different redshifts, will cause replication effects when piled in a mock light-cone, namely the regular repetition of structures in mock catalogues or images. *Transverse replications* are due to the fact that the same volume, in the same state of evolution, is used several times to fill an observing cone across the line of sight. *Radial replications* occur because the same volume, although taken at different cosmic times, is repeatedly used to fill the observing cone along the line of sight. Because the largest structures evolve slowly (i.e. over several time-steps), they will create pseudo-periodicity in mock catalogues or mock maps. In Fig. 1 (see also Fig. 2), we show how replications create an artificial perspective effect in catalogues (left hand side panel). Replication effects can be suppressed with the "random tiling" method (right hand side panel), which we describe here.

3.2.1 random tiling

Building a catalogue from the inputs described above consists of distributing the simulated galaxies in an observing cone, and computing their apparent properties in this new geometry. First, we define a three-dimensional pavement of cubic underlying boxes of side L_{box} (= 100 comoving h^{-1}Mpc). Then, we fill the underlying boxes inside the light cone with galaxies in the following way :

• determine the time-steps $i = n, ..., n+k$ which will be needed in order to fill the current underlying box, knowing that time-step i will be used to fill the light cone between $[z(i-1)+z(i)]/2$ and $[z(i)+z(i+1)]/2$;
• to each of these snapshots, apply the same transformation, which is a random combination of the following transformations :
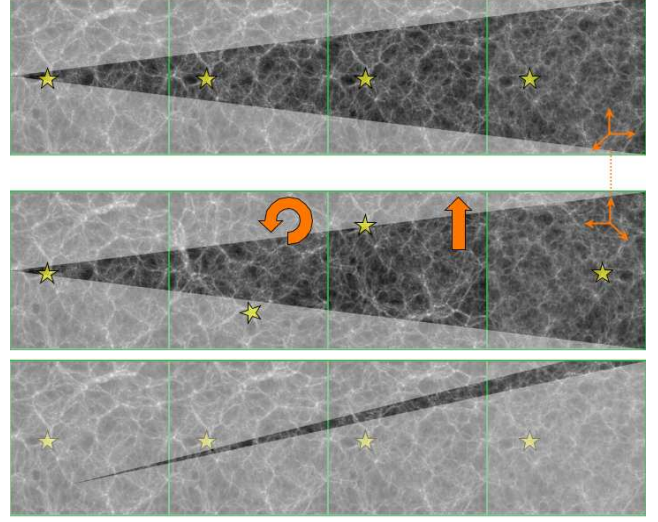

Figure 2. Illustration of the cone-making process. On the *top panel* we show the result of a straight-forward tiling of snapshots. In this case, the structures (such as the yellow star) appear repeatedly along the line of sight. The *middle panel* shows the effect of the three types of transformation we apply randomly to each underlying box. Thanks to rotations, translations, and inversions, the underlying boxes are decorrelated one from another. On the *lower panel* we show how it is possible to avoid re-shuffling underlying boxes to generate a pencil-beam type field.

– a *shift* of random amplitude (between 0 and L_{box}) in each of the three directions (x, y, z),
– a *rotation* of 0, $\pi/2$, π, or $3\pi/2$ around each axis,
– the *inversion* of one of the axes picked randomly (e.g. $x \mapsto -x$), or none;

• use the transformed positions and velocities of galaxies to include them in the light cone and compute their apparent properties;
• move on to the next underlying box, and repeat the previous steps until the light cone is filled.

The first step allows a galaxy in the cone to be taken from the output box which has the closest redshift to the galaxy's redshift relative to the observer. This has the advantage of picking galaxies at a stage of evolution as close as possible to that they would have if we had continuous outputs.

In the second step, the shifting, rotating and inverting of the underlying boxes is done to suppress replication effects. The shuffling of the underlying boxes, outlined with thick lines in figure 2, decorrelates them from one another, thus suppressing replication effects as well as any information on scales larger than the box size. Although breaking the continuity of the density field makes us loose a fraction of spatial information (see section 4.1), we chose this solution because we have good control on this information loss.

For deep pencil-beam surveys it is possible to avoid replication effects simply by choosing an appropriate line of sight so that the light-cone will intersect different regions of each underlying box. It is better in this configuration not to shuffle the boxes so as to keep all the spatial information. This is an option which is implemented in our code, and illustrated on the lower panel of Figure 2. Note however
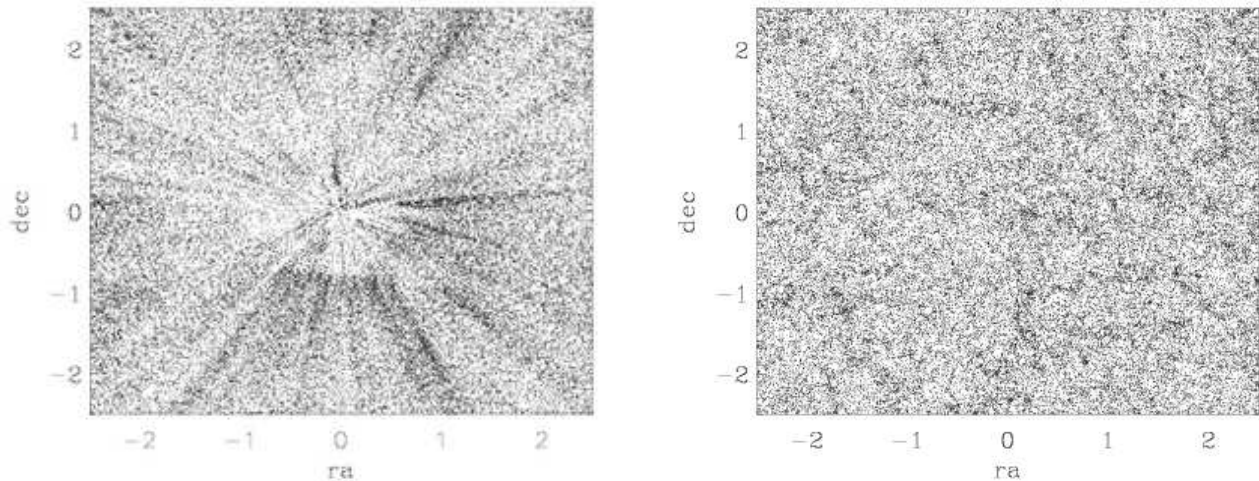
Figure 1. Angular projection of mock catalogues of 5×5 square degrees (the "r.a." & "dec." labels are arbitrary labels for two orthogonal directions on the mock sky). Each point represents a bright galaxy. The left hand side panel shows a catalogue in which the snapshots were piled together without any randomisation. This leads to replication effects, similar to a perspective effect (in an expanding/evolving universe). The right hand side panel shows a catalogue containing the same galaxies, which was made using the random tiling method. All replication effects have disappeared.

that, although the density field is continuous throughout the cone, clustering information on scales larger than L_{box} is still missing, because it is not contained in the DM simulations : replication effects can be suppressed but not finite volume effects (see Sec. 4.2).

It is also possible to chose the position of the observer, relative to the first underlying box, in both options. This allows to test for cosmic variance on local sources, and is useful to understand the statistical significance of the bright end of galaxy counts.

3.2.2 apparent magnitudes

Because we use a finite number of time-steps, (i) galaxies are picked at a cosmic time which is different from that corresponding to their distance in the mock light cone, (ii) the SEDs are not convolved with the filters at the exact redshifts, and (iii), IGM extinction is not computed for the correct redshifts. Point (i) means that an individual galaxy is not taken at the stage of evolution it would have in the case of continuous outputs. However, this does not affect the statistical properties of the mock catalogues because the overall galaxy population does not evolve much between time-steps, on average. This issue is also discussed in Sec. 4.3.1.

Points (ii) and (iii), we correct for as follows. We define corrected observer-frame absolute magnitudes as

$$M_{cor} = M[z(i)] + \frac{\mathrm{d}M}{\mathrm{d}z} \times [z(d) - z(i)], \qquad (4)$$

where $M[z(i)]$ is the observer-frame absolute magnitude computed by GALICS at redshift $z(i)$ of time-step i (Eq. 1), $z(d)$ is the redshift of the galaxy evaluated from its comoving distance d to the observer in the mock light cone and taking into account the peculiar velocity of the galaxy along the line of sight, and $\mathrm{d}M/\mathrm{d}z$ is defined in Eq. 3. Note that this derivative only accounts for distance effects (K-correction

and IGM extinction) and not evolution (point (i) above). The apparent magnitude of a galaxy is then obtained with the luminosity distance d_L :

$$m = M_{cor} + 5 \log \left(\frac{d_L}{10 \text{ pc}} \right). \qquad (5)$$

Thanks to the first order correction of magnitudes, the distribution of galaxies in apparent colour-colour plots is continuous. This is especially important for colour selections of distant galaxies as shown in Blaizot et al. (2004).

In Fig. 3, we show an example light cone with a detection limit close to that of the 2dFGRS (Colless et al. 2001). Each point represents a galaxy with $b < 19.5$, and the colours indicate apparent $B - V$ colour of the galaxies.

3.3   Mock Maps

Two types of maps are useful to address different issues :

• *pre-observation maps* are a simple projection of a mock catalogue on the sky. The only additional assumption required here is the functional form for the galaxy light profiles (e.g. an exponential disc).
• *post-observation maps* include, in addition, realistic modelling of the characteristics of the telescope/detector combination (e.g. diffraction effects, readout noise, photon shot noise). Where appropriate, atmospheric effects can also be included (e.g. seeing, air glow). SkyMaker (Erben et al. 2001) is a useful tool for producing post observation maps.

3.3.1   Pre-observation maps

Consistent with the modelling of galaxies in GALICS, we display disc components with an exponential profile, and bulges and starbursts with a Hernquist profile (Hernquist 1990, equations 32-34). The profiles are truncated at about ten
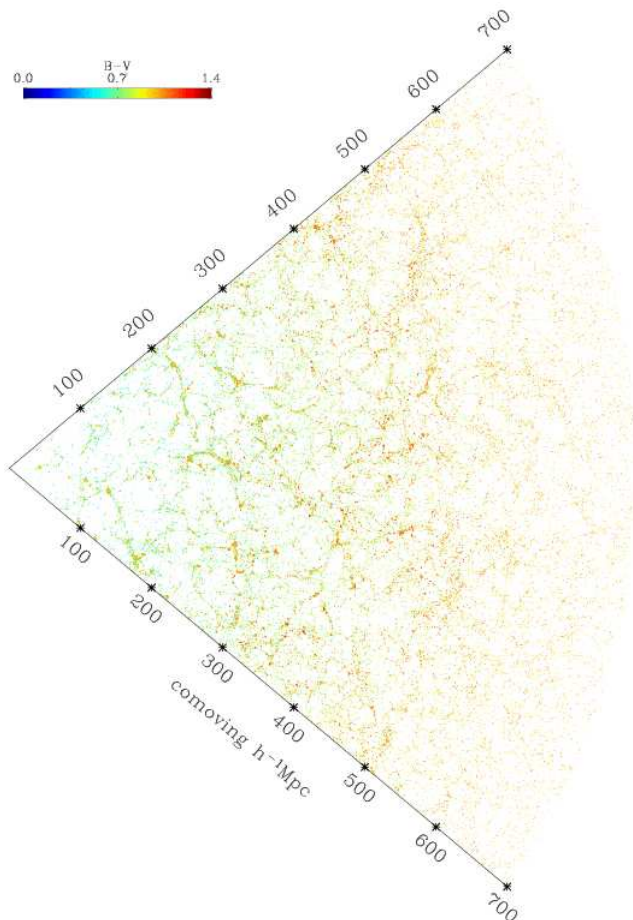
Figure 3. Mock 2dFGRS : each point represents a galaxy brighter than $b = 19.5$ from a light-cone of 75×4 square degrees. The cone was truncated at a comoving distance of $700h^{-1}$Mpc, corresponding to $z \sim 0.23$. The colours of the points indicate the apparent $B - V$ colours of the galaxies according to the colour table in the upper-left corner.

times the component's half-mass radius, and the Hernquist profile is also dimmed exponentially starting at five times the half-mass radius. To gain speed, we build a series of face-on disc and bulge templates on grids of different resolutions (e.g. from 2×2 to 2048×2048 pixels), each one normalised to unity. When adding a galaxy's contribution to the final map, we chose the template which has a resolution just above that of the final map. For bulges, as they are assumed to be spherically symmetric, the projection is straightforward, and only rescaling of the template to the component's size is required. For discs, we first have to flatten the template to account for inclination[1], and then to rotate it to map the disc's orientation. Eventually, we project the transformed templates on the final map grid, multiplying each template by the flux of the component it represents. The total flux on the final map is thus the sum of the fluxes of all galaxies in the light cone, except when they are truncated on the border of the image.

Some aliasing effects appear because of the projection

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[1] This inclination is the same as that used in GALICS to compute the extinction of light by dust.

of the tilted template grids on the final map, but these will be washed out when the map is convolved with a PSF afterwards. Since the aim of this tool is to produce pre-observation maps (with a resolution that should be higher than the final post-observation map), there is not much point in correcting this effect via bilinear interpolation or other CPU-expensive methods.

Note that the images produced this way are not limited in magnitude (up to the resolution limit of the simulation) and include the contributions of all galaxies in the cone. It is important that all sources are added (even though some may be fainter than the detection limit) for estimating the background intensity. This is particularly relevant to far infrared or sub-millimetre surveys which are limited by confusion and where the background contains a good part of the information.

Example mock maps are shown in Fig. 4. The left hand side panel shows an optical view (R band) of a 3×3 square arcmin field, and the right hand side panel shows the far IR view of the same field (at 170μm). This latter image was convolved with a Gaussian PSF of width 10 arcsec to mimic an observation by the PACS instrument on-board Herschel. No noise was added to these mock maps.

Finally, note that MOMAF allows to generate all-sky maps, using the HEALPix pixelisation (Górski et al. 2002) chosen by the Planck consortium.

3.3.2 Post-observation maps

It is considerably more difficult to generate post observing maps because separate modelling is required for each telescope/detector combination. MOMAF is designed to feed Instrument Numerical Simulators (INS) with realistic catalogues or pre-observation maps. In the optical and near-infrared domain, a ready general tool for post observing map generation is available in Skymaker (Erben et al. 2001). We briefly discuss this general INS here as an example of MOMAF possibilities.

Skymaker is an image simulation program, originally designed to assess SExtractor detection and measurement performances (Bertin & Arnouts 1996). The code (currently at version 2.3.4) has been much improved since. It is capable of simulating star and galaxy images with high level of accuracy. Galaxies in Skymaker are modelled as a combination of a de Vaucouleurs bulge and an exponential disk. Various sources of noise and convolution with the Point Spread Function can be included as desired.

There are two input files required for Skymaker to generate an image. One is a configuration file specifying the characteristics of telescope and detector and the seeing conditions. The second is the source list, which can include stars and galaxies. A typical line for a galaxy in a source list for Skymaker includes the "total" magnitude, the bulge-to-total luminosity ratio, bulge equivalent-radius in arc-second, projected bulge aspect ratio, bulge position angle in degrees, disk scale length in arc-second, disk aspect ratio and disk position angle in degrees. All of these are natural outputs of GALICS/MOMAF, as discussed in the previous section. There is no provision for adding starburst components in Skymaker. We use a workaround for such cases. We add the burst components from GALICS as additional bulges, with a scale-length obtained from GALICS, a bulge-to-total luminos-
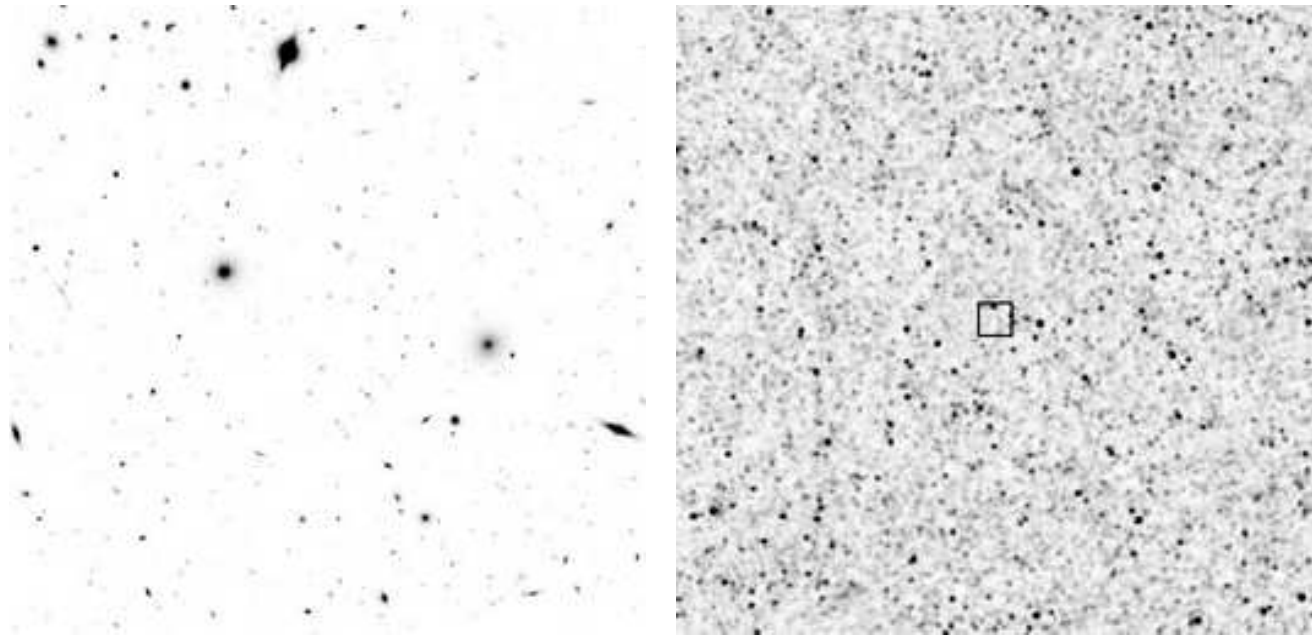
Figure 4. *Left hand side panel :* mock map of a 3×3 arcmin^2 field, in the R band. *Right hand side panel :* mock map of a 1 deg^2 field, at 170μm, convolved with a Gaussian PSF of width 10 arcsec. The central square in this image outlines the field shown in the optical in the left-hand side panel.

ity ratio of 1.0 and the appropriate starburst magnitude as the "total" magnitude.

4 LIMITATIONS OF THE METHOD

Ideally, one should build a mock catalogue from a simulated volume much larger than the light-cone, and output smoothly the physics by propagating photons towards the observer through the expanding simulated universe. Although such simulations are becoming feasible today, their computational cost is still prohibitive. The method we propose with MOMAF stems from the same philosophy as GALICS and consists of extracting as much information as possible from a given simulation, and use that to build realistic catalogues at relatively low computational expense. Of course, however sophisticated the method we use, several limitations appear in MOMAF mock catalogues because they are built from the replication of finite information. The purpose of this section is to understand how the replication process affects our predictions.

The most important limitations of MOMAF result from the fact that we use a finite volume to describe the whole Universe. In order to do so, we have to replicate the simulated volume many times along and across the line of sight. Now, because we use the random tiling method to proceed with these replications, some clustering information is lost. This results in a negative *random tiling bias* which is discussed in Sec. 4.1. A more subtle effect comes from the fact that the finite volume of the simulation used to build a mock catalogue does not describe density fluctuations on large scales. Thus these fluctuations will be missing from the mock catalogues. This results in biases on counts variance estimates and correlation function estimates. These *finite volume effects* are described in Sec. 4.2.

In this section, we also check that other possible effects are under control. In Sec. 4.3.1, we investigate the effect of finite timestep on the apparent properties of galaxies in mock catalogues. In Sec. 4.3.2, we check the impact of mass resolution of the root simulation on different observable statistics.

4.1 Random tiling bias

A negative bias on correlation functions is introduced in mock catalogues by the *random* tiling approach, which comes from the fact that we decorrelate pairs of galaxies from one underlying box to the other when re-shuffling them to suppress (periodic) replication effects. Here, we first estimate this bias on the spatial two-point correlation function, and then project the results to derive the bias on the angular correlation function.

4.1.1 Spatial correlation function (SCF)

The spatial correlation function (SCF) can computed by measuring the number of pairs of objects separated by a given distance. If one uses the estimator of Landy & Szalay (1993, hereafter LS93) :

$$\xi(r) = \frac{DD(r) - 2DR(r) + RR(r)}{RR(r)}, \tag{6}$$

where $DD(r)$, $DR(r)$ and $RR(r)$ are the number of data-data, data-random and random-random pairs separated by distance r respectively, one gets, through the logarithmic differentiation, the relative error on $\xi(r)$:

$$\frac{\delta\xi(r)}{\xi(r)} = \frac{\delta DD(r)}{DD(r)}, \tag{7}$$

because the random sample is not affected by the replication technique.

Now, take a galaxy near the edge of an underlying box (say at a distance $d < r$ from the edge). The mean number of pairs that will be missed for this galaxy, due to replication, is :

$$\widetilde{\delta DD}_g(d, r) = 2\pi r^2 n\xi(r)(1 - d/r)\mathrm{d}r, \qquad (8)$$

where n is the mean density of galaxies, and the subscript g denotes that DD is the number of pairs lost for *one* galaxy. The tilde over DD denotes that we only consider pair loss through one side of the box. To compute the loss of pairs due to one side of the underlying box, we integrate the previous equation over d, from 0 to $2r$, namely :

$$\begin{aligned} \widetilde{\delta DD}(r) &= 2\pi r^2 n\xi(r)\mathrm{d}r \int_{d=0}^{d=2r} nL_{box}^2(1 - d/r)\mathrm{d}d \\ &= 2\pi r^3 L_{box}^2 n^2\xi(r)\mathrm{d}r. \end{aligned} \qquad (9)$$

Note that we neglected corner effects here. This is justified by the fact that one should always consider separations much smaller than the size of the box (i.e. $r \ll L_{box}$).

To get the pair loss over a whole box, simply multiply the previous result by 6 (the number of sides) :

$$\delta DD(r) = 12\pi r^3 L_{box}^2 n^2\xi(r)\mathrm{d}r. \qquad (10)$$

Had we not broken the continuity in the density field between each underlying box, the number of pairs would simply be, for a whole box :

$$DD(r) = 4\pi r^2\mathrm{d}r n(1 + \xi(r)) \times nL_{box}^3, \qquad (11)$$

where the first right hand term is the number of pairs expected for one galaxy, and the second right hand term is the number of pairs in the box. The relative error in the number of pairs, due to shifting the underlying boxes is thus :

$$\frac{\delta DD(r)}{DD(r)} = 3\frac{r}{L_{box}}\frac{\xi(r)}{1 + \xi(r)}. \qquad (12)$$

For a numerical estimate, consider the Lyman break galaxies (LBG) population. For these galaxies, we expect $\xi(r) = 1$ at $r_0 \sim 6 \ h^{-1}$Mpc. Thus, for our simulation, with $L_{box} = 100 \ h^{-1}$Mpc, one finds $\delta\xi/\xi < 10\%$ for LBGs, at $6 \ h^{-1}$Mpc.

On figure 5, we show the theoretical underestimation on spatial correlation function measurements from our catalogues. In the plots, we assume a correlation of the form $\xi(r) = (r/r_0)^\gamma$, and we let r_0 and γ vary. For a wide range of these two parameters, the error due to transforming the underlying boxes is less than 10% from 1 to 10 h^{-1} Mpc.

On figure 6, we show a measure of the bias on ξ introduced by randomising the boxes. To do this, we cut our snapshot volume into 8^3 sub-boxes to which we applied translations, rotations and inversions as described above. We then measured $\xi(r)$ on the original snapshot (solid curve) and on the shuffled snapshot (dashed line). We plotted on figure 6 the prediction for the bias according to the above calculation as the dot-dashed line. The agreement between the measurement and the analytical prediction is very good on scales up to $\sim 1/5L_{box}$ (with L_{box} being $100/8h^{-1}$Mpc, as shown by the vertical line on figure 6). In the previous section, we showed that this scale is where finite volume effects come into play. Although most finite volume effects are not present here because we use all the sub-boxes to fill the simulated volume, the randomisation of sub-boxes kills any
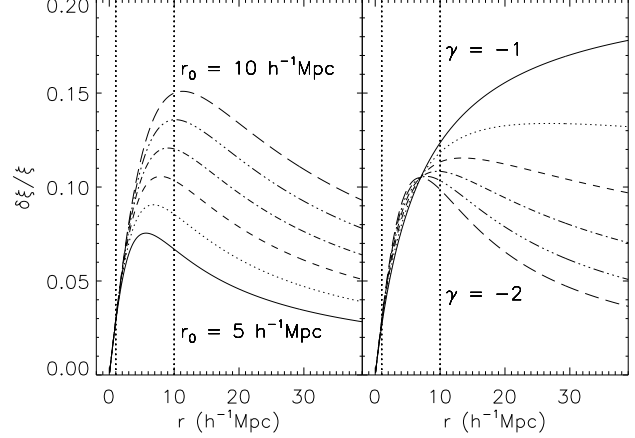

Figure 5. Expected relative underestimation on the spatial correlation function, assuming $\xi(r) = (r/r_0)^\gamma$. The left hand side panel shows variations with r_0, with $r_0 = 5, 6, 7, 8, 9, 10 \ h^{-1}$Mpc from the bottom curve to the top curve, γ being fixed to -1.8. The right hand side panel shows the dependence on γ, with γ spanning the range [-2,-1], from bottom to top, r_0 being fixed to $7h^{-1}$Mpc. On each panel, the right hand side vertical line shows the approximate upper limit of validity of measurements of ξ (one tenth of L_{box}). The left hand side vertical line roughly indicates the size of a cluster, below which our spatial information is uncertain.
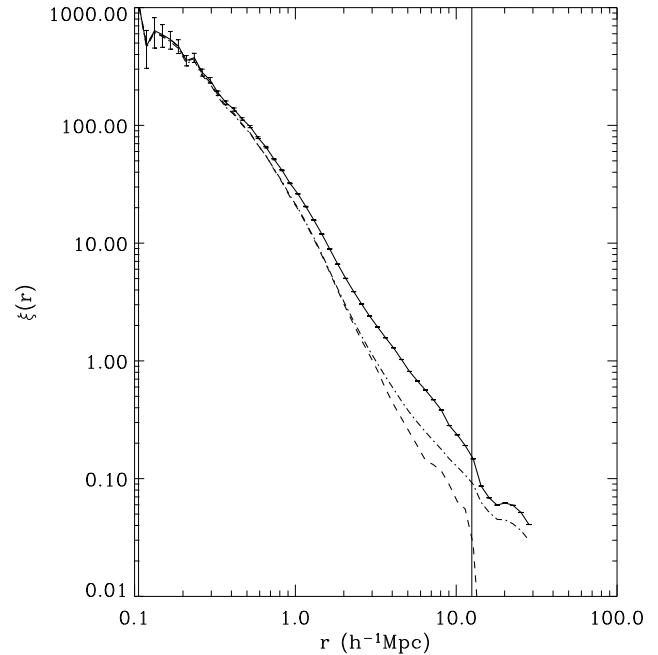

Figure 6. Measured spatial correlation function of the dark matter particles in the $z = 0$ snapshot (solid curve) and on the same snapshot re-shuffled (dashed line). Also plotted is the prediction from Eq. 12 (dot-dashed line). The error bars attached to the solid line are Poissonian errors. The vertical line indicates the size of sub-boxes used for the re-shuffling.

signal at scales larger than a sub-box. This is not described by the above analytic calculation and results in the sharp cutoff of ξ at $\sim 1/5L_{box}$.

4.1.2 Angular correlation function (ACF)

We can use the bias on the SCF derived above to evaluate that on the ACF. Let's first remember that, in the small angle approximation, the angular correlation is related to the spatial one by (Bernardeau et al. 2002),

$$w(\theta) = \int \mathrm{d}\chi \; \chi^4 D_M(\chi)\psi^2(\chi) \int_{-\infty}^{\infty} \theta\xi(r)\mathrm{d}x, \tag{13}$$

where χ is the radial distance, D_M is the angular distance ($D_M = \chi$ in a flat universe), ψ is a selection function satisfying $\int \chi^2\psi(\chi)\mathrm{d}\chi = 1$, and r is the separation distance, related to the angular separation θ and the integration variable x through the relation $r = D_M\theta(1+x^2)^{1/2}$. If we now introduce the bias on ξ as $\xi \mapsto \xi + \delta\xi$, with $\delta\xi$ given by Eq. (12), we can derive the corresponding bias on the ACF from Eq. 13 :

$$\delta w(\theta) = \int \mathrm{d}\chi \; \chi^4 D_M(\chi)\psi^2(\chi) \int_{-\infty}^{\infty} 3\theta\frac{r}{L_{box}}\frac{\xi^2(r)}{1+\xi(r)}\mathrm{d}x \tag{14}$$

Assuming that ξ can be written as the power law $(r/r_0)^{-\gamma}$, and using $D_M(\chi) = \chi$ in a flat universe, we end up with :

$$\delta w(\theta) = \frac{3r_0^\gamma\theta^{2-\gamma}}{L_{box}} \tag{15}$$
$$\times \int_0^{\infty} \mathrm{d}\chi\chi^{6-\gamma}\psi^2(\chi) \int_{-\infty}^{\infty} \frac{(1+x^2)^{1/2-\gamma}\mathrm{d}x}{\left(\frac{x\theta}{r_0}\right)^\gamma + (1+x^2)^{-\gamma/2}}$$

Note that for this result, we also assumed that the SCF does not vary with redshift. This is obviously wrong in general but is justified if the selection function ψ is narrow enough (e.g. for LBGs).

Finally, using equations (13) and (15), and deciding on a selection function, one can compute numerically the relative bias induced on ACF measurements by the transformations of underlying boxes. An example is given in figure 7.

4.2 Finite volume effects

Several limitations arise because we use a finite volume to describe the whole Universe. They are basically due to the fact that a finite volume V does not describe density fluctuations on scales typically larger than $\sim V^{1/3}$. In other words, although the mean number of galaxies in a simulation can be tuned to fit that observed in the Universe, the simulation does not describe the dispersion about this mean value. How this affects statistics from our catalogues is the question we address in this section. The simplest statistic we are interested in is galaxy counts, as a function of magnitude or redshift. Mock catalogues can be used in two ways : (i) to normalise models, and (ii) to estimate errors (including cosmic variance). In Sec. 4.2.1, we discuss how finite volume affects both the counts and their variance. Then, in Sec. 4.2.2, we describe the bias on correlation functions introduced by finite volume effects.
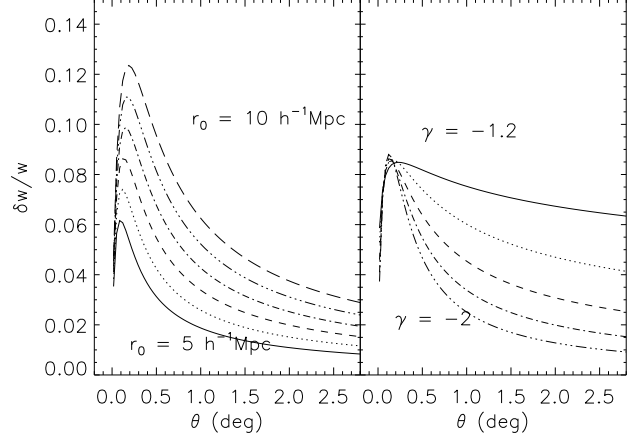

Figure 7. Expected relative underestimation on the angular correlation function, assuming $\xi(r) = (r/r_0)^\gamma$. The left hand side panel shows variations with r_0, with $r_0 = 5, 6, 7, 8, 9, 10 \; h^{-1}\mathrm{Mpc}$ from the bottom curve to the top curve, γ being fixed to -1.8. The right hand side panel shows the dependence on γ, with γ spanning the range [-2,-1], from bottom to top, r_0 being fixed to $7h^{-1}\mathrm{Mpc}$. The selection function chosen here is simply a top hat centred at $\chi = 2000h^{-1}\mathrm{Mpc}$ and of width $500h^{-1}\mathrm{Mpc}$.

4.2.1 Effects on estimates of counts variance

Two variances are relevant for counts in mock catalogues. The first is the variance which tells us about the dispersion of number counts from mock catalogues each generated from a *different* simulated volume. The second is the variance that describes the dispersion in number counts from mock catalogues made from a *unique* simulated volume. This variance tells us to what extent we can estimate cosmic variance with mock catalogues based on a given simulation. Let us proceed to virtual experiments to understand these quantities.

● Imagine we have a large number N of simulations at hand, all describing an equal volume V, but with initial conditions drawn from a much larger volume. From each simulation, we build a mock catalogue using MOMAF, and then count galaxies brighter than some magnitude limit. Finally, we measure the variance σ_1^2 of the counts obtained in this way. Now, imagine that we also have N mock catalogues, each generated using an ideal technique and a simulated volume much larger than that of the light-cone. Call σ_2^2 the variance in the counts measured from these catalogues. In the case where the volume of the cone is much smaller than volume V, the two above variances will be equal. In the more realistic case where V is smaller than the volume of the cone, one will measure that $\sigma_1 > \sigma_2$: replication enhances the bias of the simulated volume, thus dispersing more counts from catalogues.

● As a second experiment, imagine one has a unique simulated volume V as above, and builds many mock catalogues from it. These catalogues will be different from one another because of the random tiling process and because the light-cone may intersect different sections of V in different realisations. As before, measure the variance σ_3^2 of the counts from these catalogues. Again, if the simulated volume is much larger than that of the light-cone, one will measure $\sigma_3^2 \sim \sigma_2^2$. In this case, one can use mock catalogues to estimate cosmic variance. However, in the case where the volume of the

Sample	Area	Selection criteria	$< z > (z_{min} ; z_{max})$	$< i > (i_{min} ; i_{max})$	θ_{100} (deg)	N_{100}
APM	100 deg^2	$17 < B_J < 20$	0.18 (0.04 ; 0.73)	65.3 (54 ; 69)	~ 10	~ 5.5
K20	1 deg^2	$K_s < 20$	0.53 (0.02 ; 1.74)	58.0 (41 ; 69)	~ 4.5	~ 12
Counts	1 deg^2	$\chi < 5500h^{-1}$Mpc	-	-	-	-

Table 2. Geometry of mock catalogues used for comparison to data from the APM and the K20 (see text), and to make K_s-band counts. All mock catalogues have a square surface, of area given in the second column. The third column states how galaxies are selected in these catalogues (in the "Counts" case, no photometric selection is applied, but the catalogues are truncated at a comoving distance of $5500h^{-1}$Mpc from the observer). The fourth column gives the mean and span of the redshift distribution. The fifth column gives the mean and range of outputs used (output 70 is $z = 0$). The sixth column gives the angular size of the simulated volume (of side $L_b = 100h^{-1}$Mpc) at the mean redshift of the sample. The last column gives the number of radial replications needed to reach the mean redshift using the full simulated volume.

cone is larger than V, σ_3 will be found lower than σ_2, because replications do not add large-scale fluctuations. In the extreme case where the cone is very large compared to the simulated volume, σ_3 will tend to zero, because the cone encloses all the information contained in V.

In Fig. 8, we show K-band counts measured from various mock catalogues having the geometry defined in the third line of Table 2 ("counts" catalogues). The shaded area shows the locus of K20 counts from Cimatti et al. (2002), including Poissonian error bars. The filled symbols and their error bars give the mean and standard deviation for counts measured from 20 mock light-cones made using the standard simulated volume. Then we cut our root simulation into 125 sub-boxes, and made a mock catalogue out of each sub-box. The open diamonds give the mean and standard deviation of counts measured from these 125 mock light-cones. Finally the open triangles give the mean and standard deviation of counts measured from 20 mock catalogues made from a single sub-box. The upper panel of Fig. 8 compares the relative standard deviations of these three measures.

First, note that changing the size of the volume used to make mock catalogues does not change the shape of the counts, but only the amplitude. This was expected since evolution is the same in all sub-boxes. This tells us that if our root simulation is well normalised, counts from mock catalogues are not sensitive to the size of the simulated volume and can thus be used with confidence to normalise theory to observations.

Second, consider the difference between open and filled triangles in the upper panel of Fig. 8, that is, estimates of cosmic variance from cones made using one sub-box or the full simulated volume. As expected, we find that using a small box leads to an under-estimate of the variance. Table 2 tells us that the number of full boxes used to describe galaxies brighter than $K_s = 20$ is about 12 along the line of sight (up to the median redshift) and 1/5 across the line of sight at the median redshift. This situation is at the limit where we can correctly estimate cosmic variance, since the light-cone only intersects a fraction of the box volume in each underlying box. For the sub-box case, the light-cones include a full underlying box at the median redshift, and about 60 sub-boxes are replicated along the line of sight to reach this redshift. In this regime, the angular correlation function is largely under-estimated at the scale of the catalogue (because it is larger than the scale of a box), and so the estimated cosmic variance is under-estimated too.
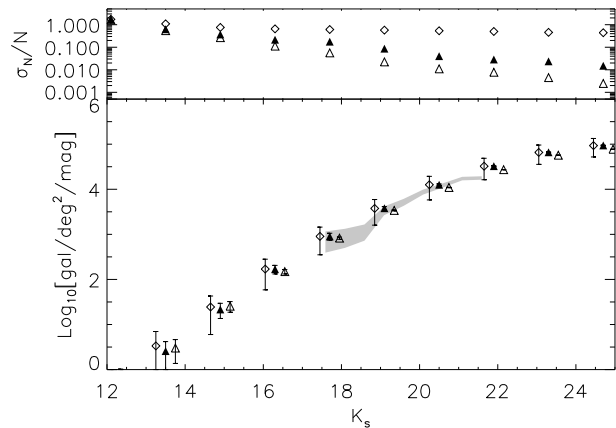

Figure 8. Effect of finite volume on K-band counts. *Lower panel*: the filled (respectively open) triangles show the mean number counts obtained from 20 catalogues made from our standard simulation (resp. from one sub-box), the error bars giving the standard deviation about this mean. The open diamonds show the mean number counts measured from 125 mock catalogues, each built from a different sub-box. The shaded area shows the locus of the K20 counts from Cimatti et al. (2002). *Upper panel*: relative standard deviation of the counts (same symbol code as lower panel).

Third, it is interesting to consider the difference between the filled triangles and the open diamonds. At bright magnitudes, the two give the same variance. This is because the volume probed by the mock catalogues is much smaller than the volume of a sub-box, so variance is well estimated with both methods, and is in fact Poissonian. At intermediate magnitudes, the volume probed by the mock cone is smaller than the full simulated volume, yet larger than that of a sub-box. Hence, the sub-box variance saturates at higher values. At the faint end, the light-cone is larger than the full simulated volume, so the variance showed with filled triangles suffers from a similar negative bias as that shown by the open triangles. The three regimes are spanned here and show that in practice, robust estimates of cosmic variance require a simulated volume much larger than the volume probed by the mock catalogue.

Redshift distributions are affected by finite volume effects in two ways. First, the variance and mean of the redshift distributions will change with box size. This effect is the same as that described above for the counts. Second, because the smaller the box, the more replications involved, repeated
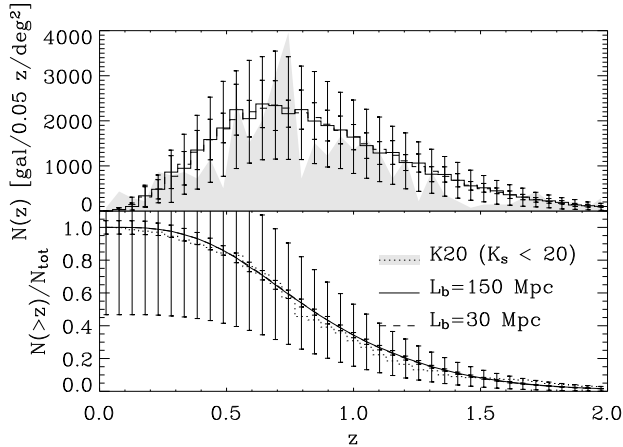
Figure 9. Effect of finite volume on the redshift distribution of a flux-limited sample. The solid (resp. dashed) curve shows the mean redshift distribution estimated from 20 mock catalogues made from the full simulated volume (resp. 125 mock catalogues each made from one different sub-box). The attached error bars give the standard deviations (the small ones correspond to the solid curve). The upper panel gives the differential distributions while the lower panel gives the normalised cumulative distributions. The shaded area and dotted line show the locus of the K20 data in the top and bottom panels (see text). The agreement with observations is very good.
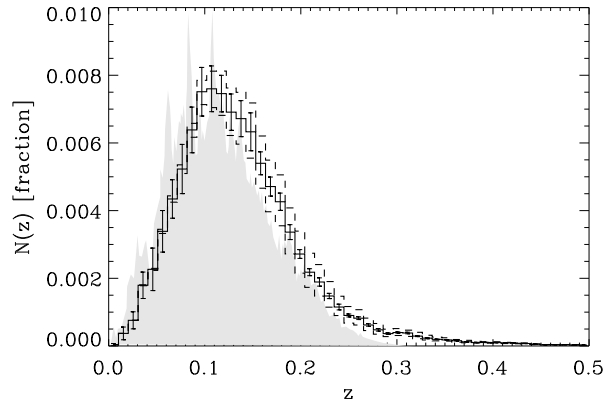

Figure 10. Comparison of GALICS to the 2dFGRS redshift distribution. The shaded shows data from Colless et al. (2001). The solid histogram (resp. lower and upper dashed histograms) shows the GALICS redshift distribution for galaxies brighter than $b_J = 19.3$ (resp. 19.2, 19.4) estimated from 20 mock surveys of 10×75 square degrees. The error bars show the dispersion in the estimates. The same arbitrary normalisation was applied to the 3 histograms.

structures may imprint periodic features in $N(z)$. Thanks to the random tiling technique, this problem is avoided. In Fig. 9, we show the differential (upper panel) and cumulative (lower panel) redshift distributions of galaxies selected as in the K20 survey (Cimatti et al. 2002). The shaded area (resp. dotted curve) indicates the locus of the data in the upper panel (resp. lower panel). In both panels, the curves (which are mostly over-imposed) show the mean distributions measured from 20 mock catalogues made with the full simulated volume and from 125 mock catalogues each made with a different sub-box. The error bars show the standard deviations about these means (the larger error bars correspond to the sub-box catalogues). Fig. 9 shows that the redshift distribution does not change in shape when the size of the box varies. Although the details of $N(z)$ will differ from one catalogue to the other, the statistical significance of redshift distributions is thus robust, and found to be in good agreement with K-band observations.

The agreement found with the K20 redshift distribution is an important success of the GALICS model, given the difficulty experienced by other models in achieving this task. The redshift distribution observed in the 2dFGRS (e.g. Colless et al. 2001) also seems to have been challenging for modellers to reproduce, and it is interesting to see how GALICS and MOMAF pass this test. The shaded area in Fig. 10 shows the redshift distribution of 2dF galaxies given by Colless et al. (2001). This distribution includes the whole survey, and thus corresponds to a *nominal* magnitude cut at $b_J = 19.45$. Because of various sources of incompleteness, however, the *effective* magnitude cut is more likely to lie around $b_J = 19.3$ (see Fig. 14 from Colless et al. 2001). The solid histogram in Fig. 10 shows the average redshift distribution measured in 20 mock surveys of 10×75 square degrees, limited in apparent magnitude at $b_J = 19.3$. The asso-

ciated error bars show the dispersion around this mean. The dashed lines show the redshift distributions corresponding to an apparent magnitude cut at $b_J = 19.2$ (lower line) and $b_J = 19.4$ (upper line). The comparison suggests that the evolution of the b_J-band luminosity function predicted by GALICS is incorrect, giving too many bright galaxies at high redshifts. Indeed, an apparent magnitude cut at $b_J \sim 19.1$ is necessary to bring our redshift distribution into better agreement with the 2dFGRS results. Let's note however that the scope of this comparison is limited in several ways. First, one should include a proper description of the complex selection function of the 2dFGRS for a more meaningful comparison. Although beyond the scope of this paper, this is readily feasible by applying the masks of the 2dFGRS to MOMAF mocks with similar geometry. Second, the dispersion showed in Fig. 10 tells us that despite the huge amount of data gathered by the 2dFGRS, cosmic variance is still quite large. Following the above discussion on finite volume effects, and looking back at Fig. 3, one sees that this dispersion is bound to be an under-estimate of the true cosmic variance because of the many replications involved in 2dFGRS-like mock surveys. This tells us that we need a bigger simulated volume to actually constrain the model : one needs to have realistic cosmic variance at a survey's size before hoping to discriminate between different models. Finally, this example shows how useful the MOMAF software is to carry out detailed comparisons of models with various datasets.

4.2.2   Effects on estimates of 2-point correlation functions

Finite volume effects alter correlation functions in a complex way. Let's first discuss what happens to the spatial correlation function (SCF) in a cubical box such as the simulated volume. The situation in mock catalogues is analogous but also includes projection effects. Following LS93 we relate the correlation function $\hat{\xi}$ contained in the simulated volume V

to the "real" ξ as

$$1 + \hat{\xi} = \frac{1 + \xi}{1 + \bar{\xi}_V}. \qquad (16)$$

At small separations (compared to $V^{1/3}$), where $\bar{\xi} \ll \xi$, the bias is negligible. At large scales, $\bar{\xi} \sim \xi$ and $\hat{\xi}$ falls down to 0. This bias directly results from the fact that the variance cannot be estimated properly at the simulated volume scale, from only one simulated volume. Let's carry out numerical tests to better understand the finite volume bias on the spatial correlation function. We again cut our standard simulation of side $L_b = 100h^{-1}$Mpc into 125 cubic sub-boxes of side $L_{sb} = 20h^{-1}$Mpc, and we measure the spatial correlation function (SCF) in all these 126 boxes, for galaxies brighter than $B = -19$. This magnitude cut leaves us with about 150 galaxies per sub-box. In Fig. 11, we plot the SCF measured from the full simulation (ξ_{100}) with diamonds, and the average of the 125 measures on sub-boxes ($\langle \xi_{20} \rangle = \sum \xi_{20}/125$) as stars. The error bars attached to the stars show the standard deviation from the 125 estimates of ξ_{20}. Comparison of ξ_{100} and $\langle \xi_{20} \rangle$ shows that finite volume effects translate into a negative bias at all scales, with a rather sharp cutoff at $r \sim L_{sb}/5$. The dashed line shows $\langle \xi_{20} \rangle$ corrected from the integral constraint given in Eq. 16. The agreement of the dashed line with the diamonds is very good at large scales. At separations smaller than $r \sim 1h^{-1}$Mpc, sub-box–to–sub-box fluctuations (both due to sparse sampling and to clustering) are responsible for the remnant discreteness bias. To understand this, let's consider the pair-weighted average of ξ_{20} :

$$\tilde{\xi}_{20} = \frac{\sum_i n_i^2 \xi_{20,i}}{\sum_i n_i^2}, \qquad (17)$$

where $\xi_{20,i}$ is the correlation function measured (with the estimator from LS93) on the n_i galaxies of sub-box i. This weighted average is shown with the triangles on Fig. 11. Notice that in the case where edge effects are negligible (i.e. at small separations), one finds

$$\tilde{\xi}_{20} \simeq \frac{\sum_i DD_i - 2DR_i + RR_i}{\sum_i RR_i} \simeq \frac{DD - 2DR + RR}{RR}, \quad (18)$$

where DD, DR, and RR are the numbers of data-data, data-random, and random-random pairs in a given separation bin for the whole simulated box. In other words, the pair-weighted average of the correlation functions of the sub-boxes is equivalent, at small scales, to using the LS93 estimator for the whole box, and is thus only affected by the integral constraint. Now, the main difference between this estimate and the estimate obtained from the full simulation (open diamonds) is that cross pairs between two sub-boxes are not regarded. In particular, as expected, this estimator converges to the biased estimator $\langle \xi_{20} \rangle$ at large scales. But at small scales, $\tilde{\xi}_{20}$ partly captures sub-box–to–sub-box fluctuations through the variations of n_i^2, and thus remains above $\langle \xi_{20} \rangle$. Still, $\tilde{\xi}_{20}$ is a combination of estimates of the SCF on small boxes, which are contaminated at all scales by the integral constraint effect. Hence $\tilde{\xi}_{20}$ remains below the "exact" result. When $\tilde{\xi}_{20}$ is corrected from the integral constraint as in Eq. 16 (solid line in Fig. 11), the result matches nearly perfectly ξ_{100}, as expected.

Fig. 12 shows how these finite volume effects affect the angular correlation function. The solid line shows the mean
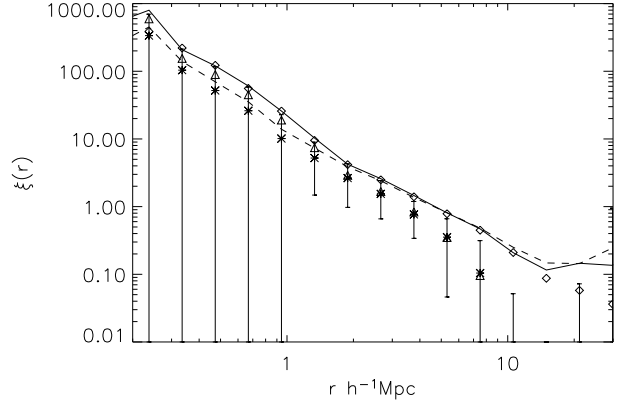

Figure 11. Spatial correlation function of galaxies brighter than $B = -19$. The diamonds show ξ_{100}, computed from the whole simulation, and the stars show the mean of estimates of ξ_{20} from 125 sub-boxes of side $L_{sb} = 20h^{-1}$Mpc (stars). The error bars show the standard deviation about this mean. The dashed line shows ξ_{20} corrected for finite volume effects (see text), and agrees with the diamonds at large separations.

of angular correlation functions measured from 20 mock APM catalogues (see Tab. 2) made from the full simulated volume, the attached error bars give the measured standard deviation. The dashed line and corresponding error bars show the mean and standard deviation of the ACF measured from the 125 mock APM catalogues made from the sub-boxes.

For the dashed line, the departure from a power-law at large scales reported in the 3-D case (see Fig. 11) occurs here at $\theta \sim 0.4$ degree. This is a direct consequence of the finite volume of the sub-box, and 0.4 degree is here about one fifth of the angular size of a sub-box at the median redshift of the survey (see Table 2). On top of this turn-around, there is an overall bias which increases slowly with separation starting at scales of about one hundredth the size of a sub-box. This is due to the projection of the bias described above for the SCF. Now, the open diamonds on Fig. 12 show the ACF measured from the APM by Maddox et al. (1996). The data are in very good agreement with our (full-box) model at scales shorter than ~ 0.1 degree – which is about $L_b/100$ at the median redshift. Long-wards of this scale, finite volume bias our ACF progressively. The comparison of our full-box $w(\theta)$ to data from the APM is similar to the above comparison between sub-box and full-box ACFs. We thus understand that the large-scale disagreement between APM and our model is not physical, but due to finite volume effects : APM data are drawn from an even larger box : the Universe !

Finally, let us come back to the issue of counts' variance in mock catalogues. Remembering that the variance of counts is basically given by the average of the angular correlation over the survey, we now clearly see how σ_3 of previous section was under-estimated. And we understand that this under-estimation will occur unless we use a simulated volume more than ten times larger than the aperture of the light-cone at the redshift of interest.
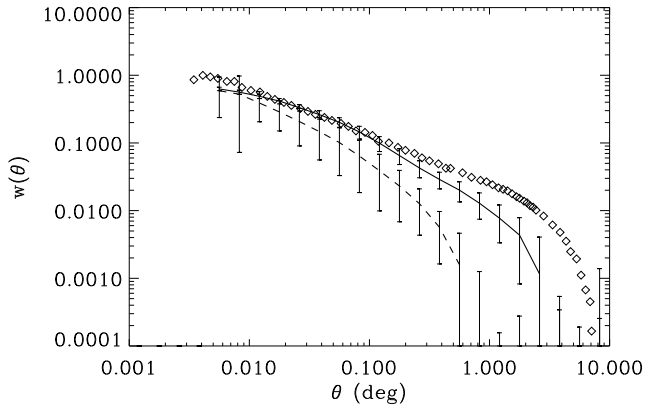
Figure 12. The solid (resp. dashed) curve show the mean ACF estimated from 20 mock catalogues made from the full simulated volume (resp. 125 mock catalogues each made from one different sub-box). The error bars give the standard deviations about these means. The cutoff in the sub-box-catalogue ACF occurs at about a fifth of the angular size of a sub-box taken at the mean redshift of the galaxy sample. Open diamonds show the ACF measured from the APM (Maddox et al. 1996).

4.3   Other effects

4.3.1   Timestep

The fact that we use a finite number of outputs, typically spaced in time by 100 Myr, could affect mock catalogues. We argued in Sec. 3.2.2 that this was not expected because even though individual galaxies can undergo a dramatic evolution during a timestep, the average properties (and their dispersion) of the overall population evolves at a much slower pace. Nevertheless, we check this hypothesis in this section by comparing statistics from mock catalogues made using different timesteps of the same simulation. Namely, we compare the counts, redshift distributions and ACFs obtained with our reference mock catalogues to those obtained with catalogues made using one snapshot out of ten[2].

The resulting counts, redshift distributions and ACFs are shown in Fig. 13, and show no significant difference between the fine and coarse time-steps. This shows that the random tiling method is robust in that the resulting mock catalogues do not depend on the time-step used in the root simulation, provided the physics was properly integrated. The fact that the properties of mock catalogues do not change with time-step shows that, at least for the selected galaxies, the K-correction and possibly a slow evolution determine the statistics. This justifies a posteriori the first-order correction made for magnitudes in Eq. 4.

The necessity of using a fine time-step to make mock observations then mainly comes from the complex analysis that can be made from them (see Blaizot et al. 2004). In this perspective, one wants to retrieve the physical properties of individual galaxies, as well as their hierarchical evolution,

---

[2] Note that in any case, the properties of the galaxies were computed using all timesteps, which is necessary in order to properly describes the physics at stake in galaxy evolution (see Hatton et al. 2003).
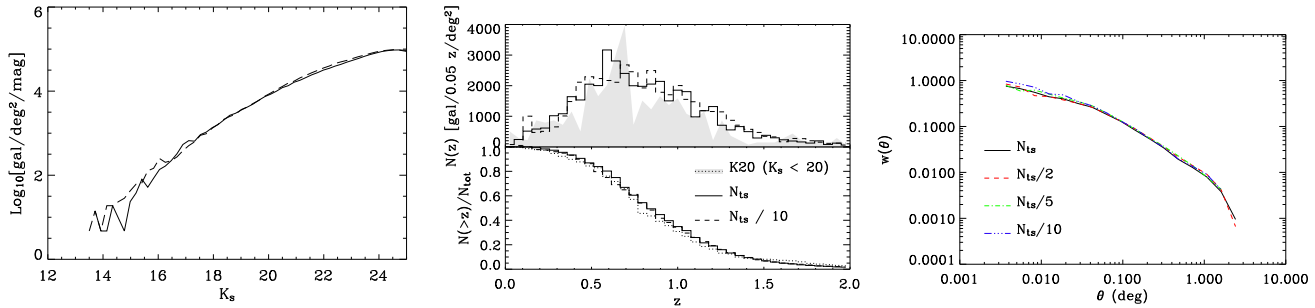
for samples selected according to observational criteria. A short time-step naturally allows to analyse evolution in more detail.

4.3.2   Mass resolution

As mentioned in Sec. 2.3, the mass resolution of the DM simulation affects galaxies in three ways : (i) incompleteness, (ii) limit redshift, and (iii) "immaturity". These three limitations, inherent to the hybrid method implemented in GALICS, will have different effects on statistics measured from mock catalogues.

(i) *Incompleteness* sets in when a fraction of galaxies of a given sample are missed because they would lie in halos below the mass resolution of the DM simulation. This effect obviously causes under-estimates of the counts at faint magnitudes. A more subtle effect is that a sample of galaxies affected by incompleteness will have a halo mass distribution biased towards high masses. Because more massive halos are more clustered, this will induce a positive bias on correlation functions. These effects cannot be corrected for except by using simulations with better mass resolution. However, as shown in Blaizot et al. (2004), it does not prevent one from using mock catalogues for studying bright galaxies, even at high redshift.

(ii) *The limit redshift* is the redshift beyond which no halo can be detected in a DM simulation. All possible galaxies at higher redshifts are thus missed by GALICS, and are hence missing from our mock catalogues. This effect, combined with incompleteness is responsible for a faint-end decrease in the counts.

(iii) *Immaturity* describes the fact that young galaxies have unrealistic properties mainly because the cooling of gas in their host haloes was not slowed down by DM accretion in sub-resolution progenitors. These galaxies only become a significant part of the overall population at redshifts higher than ~ 2 in our standard simulation, and they can be easily flagged and removed from a sample from our database. In terms of apparent magnitudes, they only significantly affect K band counts at $K > 24$.

The natural, and foreseen, solution to these limitations is to increase the resolution of root DM simulations. We come back to this perspective in the conclusions.

5   DATABASE AND WEB INTERFACE

The current implementation of the GALICS hybrid model of hierarchical galaxy formation is a package that includes three main routines (see Fig. 14). First, `HaloMaker` identifies haloes in each of the output snapshots. Second, `TreeMaker` constructs the halo merging history trees from the list of halos in all the snapshots, and computes the dark matter properties of each of the halos. Third, `GalaxyMaker` deals with the fate of baryons within the merging history trees. It computes the properties of hot gas in halos, and follows galaxy formation and evolution. The outputs are a list of properties (including absolute magnitudes in standard photometric bands) and rest-frame spectra for all galaxies in snapshots. The information produced by a given GALICS post-processing of the simulation (defined by the choice of the

Figure 13. *Left hand side panel* : number counts from a mock catalogue using all available time-steps (solid curve) and one snapshot out of 10 (dashed curve). *Middle panel* : redshift distributions using all snapshots (solid histograms) or one out of 10 (dashed histogram). The data from the K20 survey are shown by the grey area in the upper panel and by the dotted line in the lower panel. *Right hand side panel* : angular correlation functions in catalogues containing all the snapshots (solid curve), one out of 2 (dashed curve), one out of 5 (dot-dashed curve) and one out of 10 (3-dot-dashed curve).
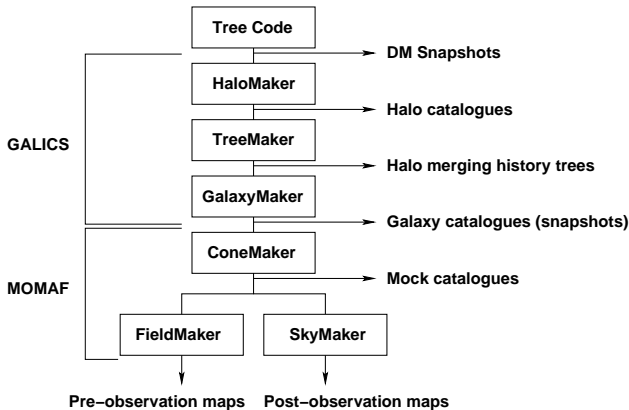

Figure 14. Organisation chart of the GALICS/MOMAF project.

astrophysical free parameters) constitutes what we hereafter call the *Archives* of this post-processing (see GALICS I). Any change in the list of input parameters will correspond to a new issue of the post-processing and new output *Archives*.

The MOMAF post-processing is also a package with three main routines, detailed in previous sections (see Fig. 14). First, `ConeMaker` generates an observing cone by integrating along the light-cone through the various snapshots, and by managing the radial and transverse replications. The output is a list of galaxies with apparent properties (including apparent magnitudes in standard photometric bands) that are computed from the *Archives*. Second, `FieldMaker` builds mock images from the observing cone, by projecting the cone galaxies onto the plane of the sky. Third, any instrument simulator can be used to transform these pre-observation images into post-observation images, or the rest-frame spectra of the *Archives* into post-observation observer-frame spectra. We have shown in the previous section how `SkyMaker` can be used in this way, because `FieldMaker` is able to generate the relevant information in the proper format. Clearly, many different cones can be generated within a single GALICS *Archive* by changing the observing point, the direction of the line-of-sight and/or the aperture. Many pre-observing images can be generated from a single cone by changing the filter response curves. And many post-observing images can be generated from a single pre-observing image by changing

the instrument simulator. The information produced by a given MOMAF post-processing constitutes what we hereafter call the *Products* of this post-processing.

At this stage, from any single simulation, we have generated a set of *Archives* and *Products* that includes tables of halo and galaxy properties, and FITS files of spectra and images. Two big issues obviously appear. First, the size of the database makes it very cumbersome. As an example, for a standard GALICS post-processing of our ΛCDM simulation, the total numbers of halos and galaxies generated in the 70 snapshots respectively amount to 1.5 and 1.8 million. The output *Archives* are about 4.5 GB for tables and 45 GB for spectra FITS files, not to speak of the *Products*. Second, the specific information that is relevant for a given user is hidden within the bulk of non-relevant information. Let's imagine for instance that we want to get the B-band absolute magnitude and total cold gas mass of a random sample of 100 galaxies brighter than apparent magnitude $I_{AB} = 20$. Extracting this information will require reading tables with many columns (the properties) and many rows (the galaxies). It may be possible to anticipate the latter issue, and generate many specific tables for many different situations and potential users, but the same pieces of information will consequently be duplicated many times, which is not the proper way to proceed.

The solution to this conundrum is well known: it consists in storing the information into a *relational database* (hereafter RDB). Here, we use the word database (with its loose meaning) for all the information we want to make available, and the word relational database (with its strict meaning) for the technical way of putting part of this information into a specific structure.

We decided to use MySQL as the relational database server. MySQL is a freely available, widely used and extremely fast database server which is capable enough for our purposes. Tools to provide Web-based access to the MySQL server are also available. The tables generated by the GALICS and MOMAF post-processing are stored in MySQL *tables*. Our database input and testing is done using several short scripts in Perl that use the Perl/DBI module. The Web front-end uses PHP4 to pass SQL queries to the MySQL database. Query outputs can either be displayed as an HTML table within the browser or down-loaded to a lo-

cal file. In this section we briefly describe the database. A quick-start guide, sample queries and descriptions of the various fields in each table are available at the GALICS web-site (`http://galics.iap.fr`).

From a single dark matter simulation, each choice in the list of the input parameters corresponds to a GALICS post-processing with its specific *Archives* and *Products*, which in its turn corresponds to a single MySQL database. The information is stored into a structure designed after the usual analysis in terms of entities, attributes and relationships, that is designed to minimise storage space and maximise query speed. Each MySQL database is consequently organised in three MySQL tables for the *Archives*, respectively called the *box*, *halo*, and *galaxy* tables, and numerous *cone* tables for the *Products*. The database scheme is illustrated in Figure 15.

(i) The box table contains general information about mean quantities at each snapshot of the simulation, such as the cosmic time, corresponding redshift, total number of halos and galaxies within the box, and integrated cosmic quantities such as the cosmic star formation rate, cold gas content, hot gas content, etc.

(ii) The halo table contains information on the halos at each time-step. Each halo is identified by a unique ID in the simulation. This information deals with dark matter (e.g. mass of the halo, virial radius, circular velocity) as well as the baryonic content of the halos (e.g. mass of hot gas and its metalicity). On top of this, we include spatial information, namely positions and velocities of the centres of mass of the halos, and hierarchical information, that is, merging history links. This information is in principle enough for one to run one's own semi-analytic model on our dark matter simulations, and thus freely test new recipes and compare results with GALICS.

(iii) The galaxy table contains the physical information we compute for galaxies: stellar masses, star formation rates, gas contents, rest-frame absolute magnitudes in a variety of filters, etc. Each galaxy is identified by a unique ID in the simulation.

(iv) The cone tables contain the positions of galaxies distributed in a mock catalogue, along with their apparent magnitudes in a variety of filters. Each galaxy is identified by a unique ID in the cone. However, because of transverse replication, different cone galaxy ID's can point to the same galaxy ID in the simulation. There are several cone tables, corresponding to different random seeds for the box shuffling process (which mimics, to some extent, cosmic variance), or to different field sizes.

Of course, the information included in the four tables is usable simultaneously in the queries, since the ID's of halos and galaxies are shared by the tables of the database and allow one to pass information from one table to another. Companion information on mock spectra and images is stored as FITS files. The rest-frame spectra are related to the galaxies in the galaxy table, whereas the images are related to a particular cone table. The observer-frame spectra are related both to galaxies and to the cone from which the galaxies are identified.

The GALICS web-site also contains a *hierarchical query* page which allows the user to retrieve hierarchical information for any galaxy in mock catalogues or snapshots. At the
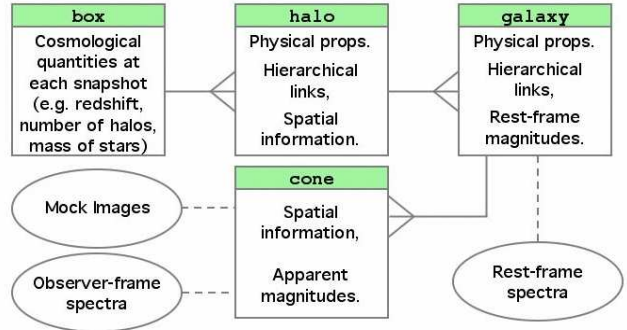

Figure 15. Data model of the GALICS/MOMAF database. There are four layers of information corresponding to four tables. The box table contains information on mean quantities computed for each time-step, e.g. redshift, cosmic time, mean SFR. The halo table contains information on the DM halos of each time-step, including their dynamical properties and merging history. The galaxy table contains information about the physical properties of galaxies at each time-step, including rest-frame magnitudes and spectra. The cone table contains information on the mock observing cone, mainly positions and apparent magnitudes. Companion FITS files include rest-frame spectra, pre- or post-observing images, and observer-frame spectra.

moment, this page contains 3 scripts. The first one allows one to view the full merging tree of a given galaxy, identified by its unique ID. The second one allow one to follow the evolution of selected properties of a galaxy along its merging history tree. Here, three options are available : (i) one can follow properties along the *main branch* which links a galaxy to its most massive progenitor at each merger; (ii) one can follow the *most massive branch*, which links the most massive progenitors at each timestep together; or (iii) one can perform a sum of additive properties on all progenitors at each timestep. Each of these options is a different way to retrieve partial information contained in the full merging history tree of a galaxy. The third script allows one to get the list of descendents or progenitors (at any redshift) of a sample of galaxies selected with any set of criteria. An example of use of this powerful script can be found in Blaizot et al. (2004). These scripts allow, for the first time, the systematic exploration of the evolution of galaxies in the framework of hierarchical galaxy formation.

6 CONCLUSIONS

In this paper, we presented the *Mock Map Facility* that takes the results of our GALICS hybrid model of hierarchical galaxy formation to make mock galaxy samples. Our method involves the construction of observing cones by integrating through the snapshots of the N-body simulation, and by using the properties of galaxies as they are computed by the GALICS post-processing. This technique builds up on the simulation and is affected by the shortcomings of the latter (mass resolution, and absence of rare objects due to the limited size of the box). It also incorporates shortcomings due to radial replication along the line-of-sight, and, for large solid angles, transverse replication. We introduced box reshuffling to minimise replication effects. The price of this technique is the loss of some signal for the correlation func-

tions (both 2D and 3D) on distances smaller than the size of the box. This loss is generally not larger than 10%. Of course, there is no signal on distances larger than the size of the box, and finite volume effects have been shown to introduce a significant (but well understood) bias on angular correlation functions.

For the purpose of analysing the limitations of our method, we compared predictions of GALICS/MOMAF to various observations. We showed that the model agrees well with K-band counts and redshift distributions. And we showed that within finite volume effects, the model also agrees well with the APM angular correlation function. These results have been obtained with the same model as used in Blaizot et al. (2004) which showed good agreement with the properties of Lyman break galaxies at $z \sim 3$. This shows that our mock catalogues can readily be used for a variety of scientific investigations.

From the mock catalogues of the observing cones, we show how to make "realistic" mock images. Since our GALICS post-processing involves multi-wavelength information from the UV to the sub-millimetre range, our mock images are produced through a wide range of standard filters. These field images can be observed through any instrument simulator. The technique is able to produce input lists for `SkyMaker`. Instrument simulators adapted to observations at infrared and sub-millimetre wavelengths can also be used.

The database produced by the GALICS and MOMAF post-processing is quite large, and has to be stored in such a way that easy access to relevant information is provided. We put the results into a relational database structure to which SQL queries can be passed through a simple Web interface. This structure has a number of well-known advantages: it optimises storage space, it makes access to the relevant information very easy, it is able to deal with simultaneous queries and updates, etc. The results of GALICS (physical properties, rest-frame magnitudes) and MOMAF (observable properties, apparent magnitudes) are stored in this database, and linked together through the standard system of a relational database model. FITS files of mock images and spectra are also available linked from the database.

The content of the database can be used for several purposes. For instance:

• comparison of mock predictions with observations through the production of a mock survey that can be processed with the same data processing pipeline as the actual survey;
• elaboration of observing strategies for forthcoming satellite missions and ground-based instruments;
• benchmark for data processing pipelines; A database populated with GALICS sources is a valuable "test set" on which to base and test the various techniques and algorithms for data reduction and analysis, for the next generation of astronomical instrumentation; the database includes the positions and magnitudes of the galaxies that are put into the mock images, and can be used as a "truth table" that has to be recovered by the data processing software;
• creation of customised galaxy samples for comparison with other models, or observational data.

We are considering improvements to this prototype database. They can develop along three axes: (i) In the mid-term future, the foreseen computer performances make repli-

cations unavoidable if a sufficient level of mass resolution for galaxy studies has to be attained. However, the improvement of the simulations will result in larger boxes that will decrease the number of radial and transverse replications, and be able to include rarer objects. (ii) The improvement of the physics within the simulations will also make better mass resolution possible, and will (hopefully) produce better results. There is no doubt also that the semi-analytic recipes have to be improved. The same cone building technique will be used also for converting the outputs of N-body simulations + hydrodynamics into mock observations. (iii) The database prototype that has been presented here will be enhanced to make it compatible with the data and metadata standards that are now being developed as part of the theoretical virtual observatory. The present MOMAF will form a valuable test-bed for testing the integration of theoretical data from simulations into the theoretical virtual observatory, which forms a part of the global Astronomical Virtual Observatory effort.

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APPENDIX A: USING THE DATABASE

In this appendix, we briefly illustrate how the database can be used to interpret observational data in the paradigm of hierarchical galaxy formation. We give four examples which exemplify the kind of information that can be exploited:

(i) synthesis of a *volume–limited sample* of galaxies, for instance at $z \simeq 0$,

(ii) synthesis of a *magnitude–limited sample* of galaxies, and related multi-wavelength information,

(iii) search for *2D and 3D spatial information* (e.g for redshift distribution, clustering), and correlation of properties with it,

(iv) search for *merging history trees* within hierarchical galaxy formation (e.g. in what type of galaxy does the material of a given high–redshift galaxy end up at $z \simeq 0$? How many progenitors does a galaxy at $z \simeq 0$ have?).

For each example, we give a typical SQL query that returns the requested subsample by querying the database. We refer the reader to the Web page (`http://galics.iap.fr/`) for additional examples, and a simple introduction to SQL syntax.

A1   Volume–limited samples

It is possible to query the database to list a series of physical properties for a subsample of galaxies with sophisticated selection criteria. As an example, select 100 galaxies at random in the $z = 0$ snapshot (that corresponds to timestep 70), with the requirement that their Johnson B–band absolute magnitude is brighter than -20, and their dispersion velocity is larger than 200 km/s. We are also interested in obtaining their absolute K–band magnitude, $B - K$ colour, morphological types and total stellar mass.

```
> SELECT gal_id, type_B2D_lum, tot_JOHNSON_B,
> tot_speed, tot_JOHNSON_K,
> tot_JOHNSON_B-tot_JOHNSON_K, tot_mstar
> FROM galaxy
> WHERE timestep=70
> AND tot_JOHNSON_B < -20
> AND tot_speed > 200
> ORDER BY RAND()
> LIMIT 100
```

The last two commands place the list of galaxies recovered by the query in random order, and then limit the output to the first 100 rows. This example takes less than 1 second to run.

A2   Magnitude–limited samples

Another type of query is to select galaxies according to their apparent magnitudes in order to mimic an observational sample or to predict what a forthcoming survey will yield. The selected sample can then be studied in the explicit cosmological context of GALICS, and the physical properties of the selected galaxies can be retrieved easily to gain insight on the nature of the "observed" objects. An example of using mock catalogues to interpret observational data is given in Blaizot et al. (2004). A crucial issue in observational galaxy formation studies is the identification of counterparts at any wavelength of galaxies observed through any given filter. It is often quite a challenge, for example, to identify the optical counterparts of far infrared sources, observed with low angular resolution. The GALICS database provides a powerful tool to address these questions as it predicts emission properties of galaxies from the UV to the sub-mm and gives the opportunity to build corresponding mock maps.

An example SQL query to retrieve a magnitude limited sample in a 1 deg^2 cone is given below, for galaxies brighter than $I_{AB} = 22.5$. We are interested in their apparent B-K colour (in the AB system), their total stellar mass, and the virial mass of their host halo. Such a query requires information that is present not only in the cone table, but also in the galaxy table and the halo table. It requires what is called a *join* in SQL syntax:

```
> SELECT cone_001.cone_id, cone_001.app_redshift,
> cone_001.JOHNSON_BAB, cone_001.JOHNSON_KAB,
> cone_001.JOHNSON_BAB-cone_001.JOHNSON_KAB,
> galaxy.tot_mstar, halo.m_vir
> FROM galaxy, cone_001, halo
> WHERE cone_001.JOHNSON_IAB < 22.5
> AND cone_001.gal_id = galaxy.gal_id
> AND halo.halo_id=galaxy.halo_id
```

This example query runs in about 30 seconds, and returns information (7 columns) for about 31000 galaxies.

The same type of selection can be used to work the other way around: one can select galaxies according to their physical properties or their dark matter halo properties, and

extract their spatial distribution and apparent magnitudes from the cone.

A3   Spatial information

Spatial information can be retrieved the same way as above, for galaxies in mock catalogues. Consider the following query :

```
> SELECT right_ascension, declination, app_redshift
> FROM cone_001
> WHERE JOHNSON_IAB < 22.5
```

This produces a table with the angular coordinates and apparent redshifts of all the mock galaxies brighter than 22.5 in the I_{AB} band within a 1 deg^2 field. There are again about 31000 such galaxies, and the query here runs in about 5 seconds.

A4   Hierarchical evolution

The GALICS project gives for the first time the opportunity to interpret observational data within the paradigm of hierarchical galaxy formation in a systematic way. One of the most important features of this theoretical framework is the notion of galaxy merging history tree. Going up or down this tree allows one to investigate the properties of the progenitors or descendents of any given galaxy at any redshift, as well as the mass build–up of that galaxy. In the galaxy table, each galaxy has a pointer towards its unique descendant at the next timestep. This minimal information is sufficient to reconstruct any merging history tree, whether forward (the list of descendants, that is a single branch as time flows) or backwards (the list of progenitors, that may be a full tree with many branches as we look back). The number of merging events for the progenitors of the galaxy whose ID is (the character string) xxyyyyyzzz is easily obtained through the following query, as well as the ID of its descendant at the next timestep (for all timesteps but the last one):

```
> SELECT nb_merge, daughter_num FROM galaxy WHERE
gal_id='xxyyyyyzzz'
```

It is necessary to run the above query recursively to build up the full merging history trees. On the GALICS website we provide PHP scripts that generate such recursive queries and pass them to the database server. Once a galaxy ID is supplied by the user through the Web interface, the ID's of its progenitors and descendants, as well as their properties, are recovered through the "recursive query" page. An interesting option allows the user to obtain the sum of the (additive) properties of all the progenitors. In such a way, the evolution of the total Star Formation Rate or the total stellar mass in all the progenitors can be easily followed. Examples of such recursive queries can be found in GALICS II and GALICS III.